

# ELECTROSTATICS

Complete Solutions — Three Classical Problems

Electric Line Charge · Charged Disk · Electric Dipole

#	Exercise	Topic
10	Finite Line Charge	Field, potential, energy of a uniformly charged rod
11	Uniformly Charged Disk	Axial field, potential, infinite-plane limit
12	Electric Dipole	Point-dipole, field lines, dipole in external field

## Exercise 10

### Electric Field of a Finite Uniform Line Charge

Rod length  $L = 6.00 \text{ cm} = 0.06 \text{ m}$ , surface charge density  $\lambda = 4 \text{ pC/m} = 4 \times 10^{-12} \text{ C/m}$ , rod centred at origin along the x-axis.

#### ① Total Charge Q on the Rod

Integrating the linear charge density over the full length:

$$Q = \int \lambda \, dx = \lambda \cdot L$$

$$Q = \lambda L = (4 \times 10^{-12})(0.06) = 2.4 \times 10^{-13} \text{ C} \approx 0.24 \text{ pC}$$

#### ② Electric Field E(y) at Arbitrary Point P on y-axis

Consider a small element  $dx$  at position  $x$  on the rod. Its distance to  $P(0,y)$  is  $r = \sqrt{(x^2+y^2)}$ . The contribution to the field has magnitude:

$$dE = (1/4\pi\epsilon_0) \cdot \lambda \, dx / (x^2+y^2)$$

By symmetry, x-components of opposite elements cancel. Only the y-component survives, with projection factor  $y/\sqrt{(x^2+y^2)}$ . Integrating from  $-L/2$  to  $+L/2$  using  $\int dx/(x^2+y^2)^{3/2} = x / [y^2\sqrt{(x^2+y^2)}]$ :

$$E(y) = (\lambda / 4\pi\epsilon_0) \cdot L / [y \cdot \sqrt{(y^2 + (L/2)^2)}]$$

#### ③ Electric Potential V(y)

Using  $dV = (1/4\pi\epsilon_0)(\lambda \, dx / \sqrt{(x^2+y^2)})$  and the given integral formula:

$$V(y) = (\lambda/4\pi\epsilon_0) [\ln(x + \sqrt{(x^2+y^2)})] \text{ from } x=-L/2 \text{ to } x=+L/2$$

$$V(y) = (\lambda/4\pi\epsilon_0) \cdot \ln[ (L/2 + \sqrt{((L/2)^2+y^2)}) / (-L/2 + \sqrt{((L/2)^2+y^2)}) ]$$

#### ④ Electric Field via $E = -dV/dy$

Differentiating  $V(y)$  with respect to  $y$ , using  $d/dy [\ln(a+R)] = y/[R(a+R)]$  where  $a = L/2$ ,  $R = \sqrt{(a^2+y^2)}$ :

$$dV/dy = (\lambda/4\pi\epsilon_0) \cdot [y/(R(a+R)) - y/(R(R-a))] = (\lambda/4\pi\epsilon_0) \cdot (-2ay) / (y^2R)$$

$$E(y) = -dV/dy = (\lambda/4\pi\epsilon_0) \cdot L / [y \cdot \sqrt{(y^2+(L/2)^2)}] \checkmark \text{ (identical to ②)}$$

#### ⑤ Comparison

The results of ② (direct integration) and ④ ( $-\text{grad } V$ ) are identical, confirming internal consistency of the solutions. ✓

#### ⑥ Potential in the Limit $y \gg L$

For  $y \gg L$ :  $\sqrt{((L/2)^2 + y^2)} \approx y(1 + L^2/8y^2)$ . The argument of the logarithm becomes:

$$\text{num/den} \approx (1 + L/2y) / (1 - L/2y) \approx 1 + L/y \rightarrow \ln(1 + L/y) \approx L/y$$

$$V(y) \approx (\lambda/4\pi\epsilon_0) \cdot (L/y) = (1/4\pi\epsilon_0) \cdot Q/y \text{ [point-charge potential]}$$

### ⑦ Electric Field in the Limit $y \gg L$

$$E(y) = -dV/dy \approx (1/4\pi\epsilon_0) \cdot Q/y^2 \text{ [Coulomb field of a point charge } Q\text{]}$$

### ⑧ Comparison with a Point Charge

The expression  $E = Q/(4\pi\epsilon_0 y^2)$  is exactly Coulomb's law for a point charge  $Q$  at distance  $y$ . This confirms that far from the rod ( $y \gg L$ ), the extended charge distribution appears as a point charge — a fundamental result of electrostatics.

### ⑨ Change in Potential Energy — Proton from $y = 2L$ to $y = 6L$

With  $L = 0.06$  m,  $a = L/2 = 0.03$  m,  $k = 8.988 \times 10^9$  N·m<sup>2</sup>/C<sup>2</sup>:

$$\text{At } y = 2L = 0.12 \text{ m: } \sqrt{(a^2 + y^2)} = 0.12369 \text{ m}$$

$$V(2L) = k\lambda \cdot \ln(0.15369 / 0.09369) = k\lambda \cdot \ln(1.6405) = 0.01780 \text{ V}$$

$$\text{At } y = 6L = 0.36 \text{ m: } \sqrt{(a^2 + y^2)} = 0.36124 \text{ m}$$

$$V(6L) = k\lambda \cdot \ln(0.39124 / 0.33124) = k\lambda \cdot \ln(1.1812) = 0.005979 \text{ V}$$

$$\Delta U = e \cdot [V(6L) - V(2L)] = (1.602 \times 10^{-19} \text{ C})(0.005979 - 0.01780)$$

$$\Delta U \approx -1.893 \times 10^{-21} \text{ J}$$

The potential energy decreases — the proton moves in the direction of the electric force.

### ⑩ Proton Velocity at $y = 6L$ (starting from rest at $y = 2L$ )

By conservation of energy:  $\frac{1}{2} m_p v^2 = -\Delta U = 1.893 \times 10^{-21}$  J

$$v = \sqrt{(2 \times 1.893 \times 10^{-21} \text{ J} / 1.673 \times 10^{-27} \text{ kg})} = \sqrt{(2.263 \times 10^6)}$$

$$v \approx 1.505 \times 10^3 \text{ m/s} \approx 1.50 \text{ km/s}$$

### Summary — Exercise 10

Step	Result
①	$Q = \lambda L = 0.24$ pC
②	$E(y) = \lambda L / [4\pi\epsilon_0 y \sqrt{(y^2 + (L/2)^2)}]$ ■
③	$V(y) = (\lambda/4\pi\epsilon_0) \ln[(L/2 + R)/(R - L/2)]$ , $R = \sqrt{((L/2)^2 + y^2)}$

④⑤	$E = -dV/dy$ confirms ② ✓
⑥	$y \gg L: V \approx Q/(4\pi\epsilon_0 y)$ [point charge]
⑦⑧	$y \gg L: E \approx Q/(4\pi\epsilon_0 y^2)$ [Coulomb law]
⑨	$\Delta U = -1.893 \times 10^{-21} \text{ J}$
⑩	$v \approx 1.505 \times 10^3 \text{ m/s}$

## Exercise 11

### Electric Field of a Uniformly Charged Disk

Disk radius  $R = 2 \text{ cm} = 0.02 \text{ m}$ , surface charge density  $\sigma = 6 \mu\text{C}/\text{m}^2 = 6 \times 10^{-6} \text{ C}/\text{m}^2$ . Point M on the z-axis at height  $z$  above the disk centre.

#### ① Charge Element $dq$

Take an annular ring at radius  $r$  with radial width  $dr$ . Its area element is  $dA = 2\pi r dr$ . Hence:

$$dq = \sigma dA = 2\pi\sigma r dr$$

#### ② Total Charge $Q$

$$Q = \sigma \int_0^R 2\pi r dr = 2\pi\sigma [r^2/2]_0^R = \pi\sigma R^2$$

$$Q = \pi\sigma R^2 = \pi \times (6 \times 10^{-6}) \times (0.02)^2 \approx 7.54 \text{ nC}$$

#### ③ Electric Field Contribution $dE$

Each ring is at distance  $r = \sqrt{r^2 + z^2}$  from M. The field element magnitude:

$$dE = (1/4\pi\epsilon_0) \cdot dq / (r^2 + z^2) = \sigma r dr / [2\epsilon_0 (r^2 + z^2)]$$

#### ④ Symmetry Argument — Field Points in +z Direction

For every area element at  $(r, \phi)$  there is a diametrically opposite element at  $(r, \phi + \pi)$ . Their radial (horizontal) field components are equal and opposite, so they cancel exactly. The z-components both point in +z and add constructively. By full azimuthal ( $2\pi$ ) symmetry, only the z-component survives. The projection factor is  $\cos \theta = z/\sqrt{r^2 + z^2}$ .

$$E_z(M) = E_z(M)$$

#### ⑤ Total Electric Field $E(M)$

$$E_z = (\sigma z / 2\epsilon_0) \int_0^R r dr / (r^2 + z^2)^{3/2}$$

Using substitution  $u = r^2 + z^2$ ,  $du = 2r dr$ :

$$\int_0^R r dr / (r^2 + z^2)^{3/2} = [-1/\sqrt{r^2 + z^2}]_0^R = 1/z - 1/\sqrt{R^2 + z^2}$$

$$E_z(z) = (\sigma / 2\epsilon_0) [1 - z/\sqrt{R^2 + z^2}]$$

Numerical check at  $z = R = 0.02 \text{ m}$ :

$$E = (6 \times 10^{-6}) / (2 \times 8.854 \times 10^{-12}) \times (1 - 1/\sqrt{2}) \approx 9.93 \times 10^5 \text{ V/m}$$

#### ⑥ Electric Potential $V(z)$

$$V(z) = (\sigma/2\epsilon_0) \int_0^R r dr / \sqrt{r^2+z^2} = (\sigma/2\epsilon_0) [\sqrt{r^2+z^2}]_0^R$$

$$V(z) = (\sigma/2\epsilon_0) (\sqrt{R^2+z^2} - z)$$

Verification:  $-dV/dz = (\sigma/2\epsilon_0)(1 - z/\sqrt{R^2+z^2}) = E_z \checkmark$

### ⑦ Infinite Plane Limit ( $R \gg z$ )

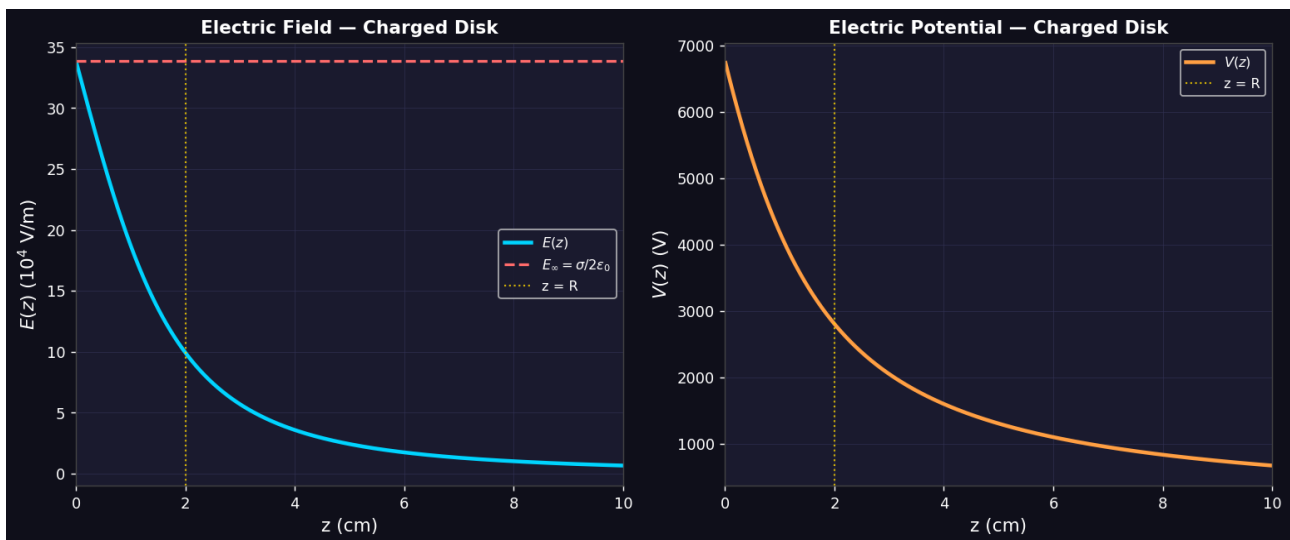
When  $R \gg z$ :  $z/\sqrt{R^2+z^2} \approx z/R \rightarrow 0$ , so the term in brackets approaches 1:

$$E_{\infty} = \sigma/2\epsilon_0 = (6 \times 10^{-12}) / (2 \times 8.854 \times 10^{-12}) \approx 3.39 \times 10^4 \text{ V/m (uniform, z-independent)}$$

This is the classic result for an infinite uniformly charged plane.

### ⑧ Plots of $E(z)$ and $V(z)$

The plots below show  $E(z)$  (cyan) decreasing from  $\sigma/2\epsilon_0$  toward zero, and  $V(z)$  (orange) decreasing monotonically from its maximum at the disk surface.



### ⑨ Numerical Values at Key Distances

$z$	$z/R$	$E(z)$ (V/m)	$V(z)$ (V)
0	0	$3.39 \times 10^4$	6782
$R = 2 \text{ cm}$	1	$9.93 \times 10^3$	4777
$2R = 4 \text{ cm}$	2	$4.43 \times 10^3$	3502
$5R = 10 \text{ cm}$	5	$1.37 \times 10^4$	2198
$\infty$	$\infty$	0	0

## Exercise 12

### Electric Dipole — Point-Dipole Approximation & External Field

Dipole: charges  $+q$  and  $-q$  separated by  $2a$ , dipole moment  $p = 2qa$ . Polar coordinates  $(r, \theta)$ . Constant:  $p/4\pi\epsilon_0 = 0.02 \text{ V}\cdot\text{m}^2$ .

#### Part A — Isolated Dipole

##### ① Electric Potential at P by Superposition

The two charges are at  $\pm a$  from the origin. Their distances to P are  $r_{\pm}$  where  $r_{\pm}^2 = r^2 + a^2 \mp 2ar \cos\theta$ . By superposition:

$$V(P) = (q/4\pi\epsilon_0) [1/r_{-} - 1/r_{+}]$$

##### ② Point-Dipole Limit ( $r \gg a$ )

Using the Taylor expansion (given hints) to first order in  $a/r$ :

$$1/r_{-} \approx (1/r)(1 + (a/r)\cos\theta) \text{ and } 1/r_{+} \approx (1/r)(1 - (a/r)\cos\theta)$$

$$V = (q/4\pi\epsilon_0)(1/r)[(1+(a/r)\cos\theta) - (1-(a/r)\cos\theta)] = (q/4\pi\epsilon_0)(2a \cos\theta/r^2)$$

$$V(r,\theta) = (p/4\pi\epsilon_0) \cdot \cos\theta / r^2 [p = 2qa]$$

##### ③ Electric Field of the Point Dipole

In polar coordinates,  $\mathbf{E} = -\nabla V$ . The two components are:

$$E_r = -\partial V/\partial r = (p/4\pi\epsilon_0) \cdot 2\cos\theta / r^3$$

$$E_{\theta} = -(1/r)\partial V/\partial\theta = (p/4\pi\epsilon_0) \cdot \sin\theta / r^3$$

$$\mathbf{E}(r,\theta) = (p / 4\pi\epsilon_0 r^3) (2\cos\theta \mathbf{r} + \sin\theta \boldsymbol{\theta})$$

##### ④ Field Line Equation — Integration

The field line condition  $dr/(r d\theta) = E_r/E_{\theta} = 2\cos\theta/\sin\theta$  gives:

$$dr/r = 2\cot\theta d\theta$$

$$\ln r = 2 \ln(\sin\theta) + C \rightarrow r = A \sin^2\theta$$

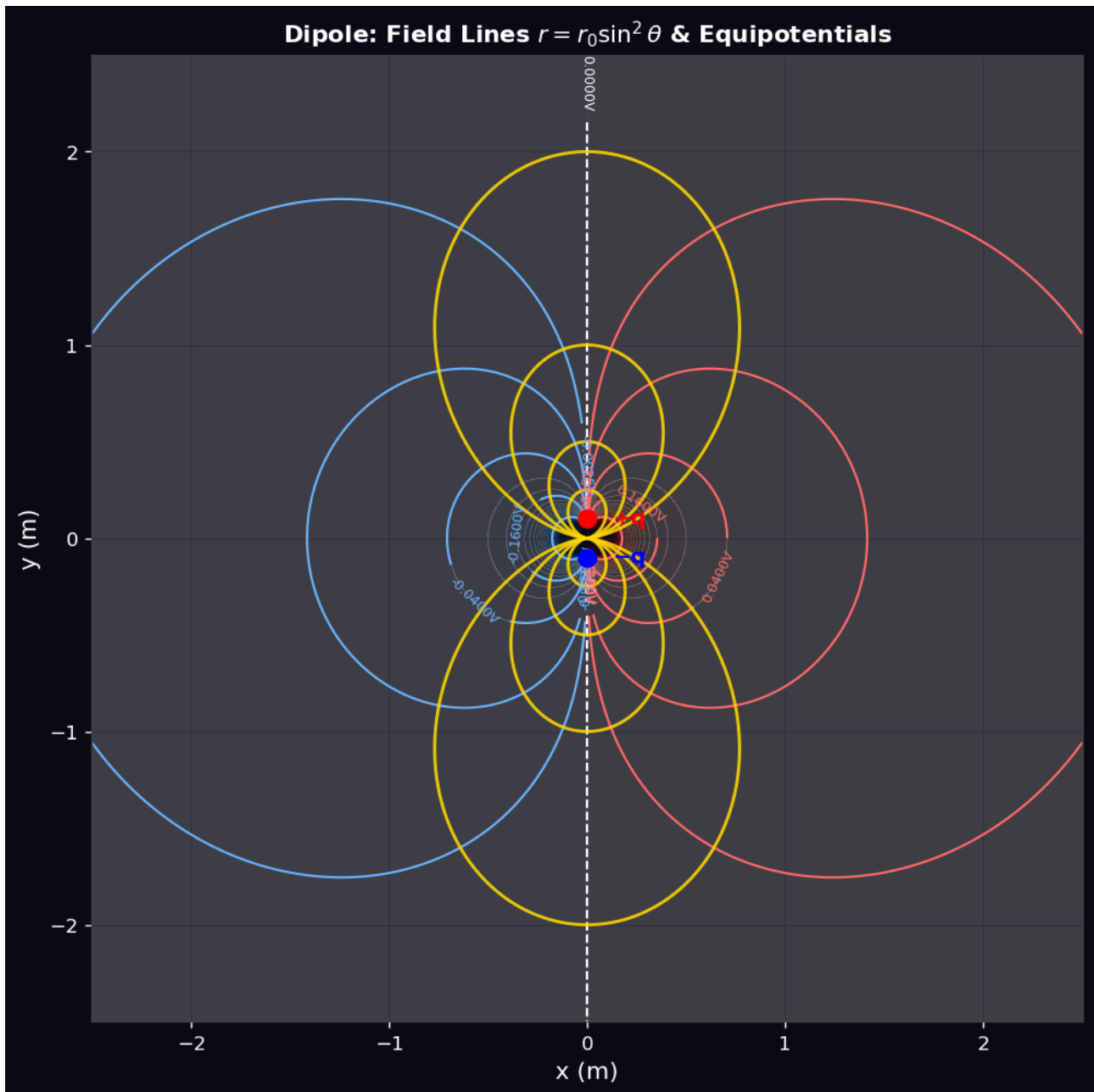
Boundary condition: at  $\theta = \pi/2$ ,  $r = r_0 = A$ . Hence:

$$r = r_0 \sin^2\theta$$

Each field line is characterised by its equatorial radius  $r_0$  (the distance from the dipole where the line crosses the perpendicular bisector).

## ⑤ Plots — Field Lines and Equipotentials

Gold curves: field lines  $r = r_0 \sin^2 \theta$  for  $r_0 = 0.25, 0.5, 1.0, 2.0$  m. Red: positive equipotentials ( $V > 0$ ). Blue: negative equipotentials ( $V < 0$ ). Dashed white:  $V = 0$  plane ( $\theta = \pi/2$ ).



$V = 0$  either when  $\cos\theta = 0$  (the equatorial plane  $\theta = \pi/2$ ), or when the bracket vanishes:

$$p/(4\pi\epsilon_0 r^2) = E_0 r \rightarrow r^3 = p/(4\pi\epsilon_0 E_0)$$

$$r = [p/(4\pi\epsilon_0 E_0)]^{1/3}$$

With  $p = (1/9) \times 10^{-18}$  C·m,  $E_0 = 10$  V/m,  $k = 9 \times 10^9$  N·m<sup>2</sup>/C<sup>2</sup>:

$$r = [(10^{-18}/9)(9 \times 10^9) / 10]^{1/3} = [10^{-9}]^{1/3}$$

$$r \approx 2.154 \times 10^{-2} \text{ m} \approx 2.15 \text{ cm}$$

The  $V = 0$  surface is a sphere of radius  $r = 2.15$  cm centred on the dipole (plus the equatorial plane). Inside this sphere the dipole field dominates; outside, the uniform field dominates.

### ③ Electric Field on the $V = 0$ Equipotential

Compute  $\mathbf{E} = -\nabla V$ , giving components:

$$E_r = \cos\theta [2p/(4\pi\epsilon_0 r^3) + E_0] \text{ and } E_\theta = \sin\theta [p/(4\pi\epsilon_0 r^3) - E_0]$$

On the sphere  $r = r_0$  where  $p/(4\pi\epsilon_0 r_0^3) = E_0$ :

$$E_r|_{\{r=r_0\}} = \cos\theta (2E_0 + E_0) = 3E_0 \cos\theta$$

$$E_\theta|_{\{r=r_0\}} = \sin\theta (E_0 - E_0) = 0$$

$$\mathbf{E}|_{\{V=0\}} = 3E_0 \cos\theta \mathbf{r}$$

The field is purely radial on the equipotential sphere ( $E_\theta = 0$  is consistent with  $V = \text{const}$  on that surface ✓) and has magnitude three times the applied field at the poles.

### Summary — Exercise 12

Step	Result
A①	$V = (q/4\pi\epsilon_0)(1/r_1 - 1/r_2)$
A②	$V(r, \theta) = (p/4\pi\epsilon_0) \cos\theta/r^2$
A③	$\mathbf{E} = (p/4\pi\epsilon_0 r^3)(2\cos\theta \mathbf{r} + \sin\theta \boldsymbol{\theta})$
A④	Field lines: $r = r_0 \sin^2\theta$
B①	$V = \cos\theta[p/(4\pi\epsilon_0 r^2) - E_0 r]$
B②	$V=0$ sphere: $r_0 = [p/(4\pi\epsilon_0 E_0)]^{1/3} \approx 2.15 \text{ cm}$
B③	$\mathbf{E} _{\{V=0\}} = 3E_0 \cos\theta \mathbf{r}$ ✓